

Stability Analysis for Transmission Problem in Thermoelasticity

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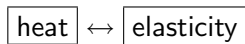
1 Introduction

- Thermoelasticity
- Transmission Problem and Its Stability

2 Transmission Problem with Concentrated Mass

- Model
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- Stability Analysis

1. Introduction: Thermoelasticity



Dynamical behavior: 1-d linear form

$$\begin{cases} u_{tt}(x, t) - \alpha u_{xx}(x, t) + \beta \theta_x(x, t) = 0, & t > 0, x \in (0, \ell), \\ \theta_t(x, t) + \gamma q_x(x, t) + \delta u_{tx}(x, t) = 0, & t > 0, x \in (0, \ell), \\ \tau q_t(x, t) + q(x, t) + k \theta_x(x, t) = 0, & t > 0, x \in (0, \ell), \end{cases} \quad (1)$$

where $u(x, t)$ is the displacement at time t , position x ; $\theta(x, t)$ is the temperature difference to a fixed reference temperature and $q(x, t)$ is the heat flux. Assume that the parameters $\alpha, \beta, \gamma, \delta, k$ are all positive parameters and $\tau \geq 0$.

the pure propagation of heat

$$\theta_t(x, t) + \gamma q_x(x, t) = 0$$

$$\tau q_t(x, t) + q(x, t) + k\theta_x(x, t) = 0$$

- $\tau = 0$: Fourier law (classical thermoelasticity)
- $\tau > 0$: Cattaneo Law (thermoelasticity of Lord-Shulman type)

The pure heat conduction are all exponentially stable to zero under these two laws.

When $\tau = 0$, system (1) is a hyperbolic-parabolic coupled system for u and θ .

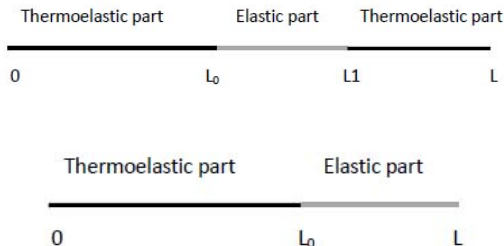
When $\tau > 0$, system (1) becomes a hyperbolic-hyperbolic coupled one.

Many kinds of thermoelastic systems have been considered, please see for example:

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Transmission Problem in Thermoelasticity



- pure elastic body \rightarrow conservative system;
- thermoelastic materials \rightarrow damping mechanism

If the thermal damping mechanism is effective in the whole domain, it always leads to exponential stability for one-dimensional linear systems; While what will happen if it is only effective in a sub-domain? The corresponding mathematical model is called a **transmission problem**.

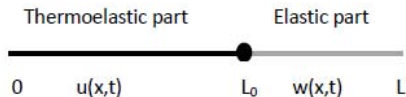
stability

The stability of various transmission problems on thermoelasticity have been considered. For **one-dimensional system**, the systems are proved to be **exponentially stable, no matter how small the thermoelastic parts are**. See for example: Racke, Rivera et al [1][2][3].

Remark: For the system with nonlinear term or the system with memory, some conditions on nonlinear term or relaxation function must be added to ensure the stability of the whole system.

- [1] A. Marzocchi, J.E. Munoz Rivera and M.G. Naso, Asymptotic behavior and exponential stability for a transmission problem in thermoelasticity, *Math. Meth. Appl. Sci.*, 2002.
- [2] L.H. Fatori, E. Lueders and J.E. Munoz Rivera, Transmission problem for hyperbolic thermoelastic systems, *J. Thermal Stresses*, 2003.
- [3] H. D. Fernandez Sare, J. E. Munoz Rivera and R. Racke, Stability for a transmission problem in thermoelasticity with second sound

2. Transmission problem with concentrated mass



$$\begin{cases} u_{tt}(x, t) - \alpha u_{xx}(x, t) + \beta \theta_x(x, t) = 0, & x \in (0, L_0), & 0 < L_0 < L, & t > 0, \\ \theta_t(x, t) + \gamma q_x(x, t) + \delta u_{tx}(x, t) = 0, & x \in (0, L_0), & t > 0, \\ \tau q_t(x, t) + q(x, t) + k \theta_x(x, t) = 0, & x \in (0, L_0), & t > 0, \\ w_{tt}(x, t) - b w_{xx}(x, t) = 0, & x \in (L_0, L), & t > 0, \end{cases} \quad (2)$$

where $u(x, t)$, $w(x, t)$ denote the displacements in intervals $(0, L_0)$ and (L_0, L) , respectively. $\theta(x, t)$ and $q(x, t)$ for $x \in (0, L_0)$ denote the temperature difference to a fixed reference temperature, and the heat flux.

We mainly consider the case $\tau > 0$, that is, the thermoelasticity of Lord-Shulman type.

Boundary and Transmission Condition

Boundary conditions at $x = 0, L$:

$$\begin{cases} u_x(0, t) = 0, & t > 0, \\ \theta(0, t) = 0, & t > 0, \\ w(L, t) = 0, & t > 0, \end{cases} \quad (3)$$

Transmission conditions at $x = L_0$:

$$\begin{cases} q(L_0, t) = 0, \\ u(L_0, t) = w(L_0, t), \\ -Mu_{tt}(L_0, t) = \alpha u_x(L_0, t) - \beta \theta(L_0, t) - bw_x(L_0, t), & M > 0. \end{cases} \quad (4)$$

The initial conditions:

$$\begin{cases} u(., 0) = u_0, & u_t(., 0) = u_1, & \theta(., 0) = \theta_0, & q(., 0) = q_0, & \text{in } (0, L_0), \\ w(., 0) = w_0, & w_t(., 0) = w_1, & & & \text{in } (L_0, L). \end{cases} \quad (5)$$

Semigroup Setting

Set

$$V^k[a, b] := \{u \in H^k[a, b] \mid u(a) = 0\},$$

$$\tilde{V}^k[a, b] := \{u \in H^k[a, b] \mid u(b) = 0\}.$$

State space: $\mathcal{H} := \left\{ (u, v, w, z, \theta, q, p)^T \in$

$$H^1[0, L_0] \times L^2(0, L_0) \times \tilde{V}^1[L_0, L] \times L^2(L_0, L) \times L^2(0, L_0) \times L^2(0, L_0) \times \mathbb{C} \mid u(L_0) = w(L_0) \right\},$$

equipped with an inner product, for

$$W_j = (u_j, v_j, w_j, z_j, \theta_j, q_j, p_j) \in \mathcal{H}, \quad j = 1, 2,$$

$$\begin{aligned} (W_1, W_2)_{\mathcal{H}} &= \int_0^{L_0} [v_1(x)\overline{v_2(x)} + \alpha u_{1,x}(x)\overline{u_{2,x}(x)} + \frac{\beta}{\delta} \theta_1(x)\overline{\theta_2(x)} \\ &\quad + \frac{\tau\gamma\beta}{k\delta} q_1(x)\overline{q_2(x)}] dx + \int_{L_0}^L [z_1(x)\overline{z_2(x)} + b w_{1,x}(x)\overline{w_{2,x}(x)}] dx \\ &\quad + M p_1 \overline{p_2}. \end{aligned}$$

System operator

$$\mathcal{A} \begin{bmatrix} u \\ v \\ w \\ z \\ \theta \\ q \\ p \end{bmatrix} = \begin{pmatrix} v \\ \alpha u_{xx} - \beta \theta_x \\ z \\ bw_{xx} \\ -\gamma q_x - \delta v_x \\ -\frac{1}{\tau}(q + k\theta_x) \\ -\frac{1}{M}[\alpha u_x(L_0) - bw_x(L_0) - \beta \theta(L_0)] \end{pmatrix}$$

with domain

$$\mathcal{D}(\mathcal{A}) = \left\{ (u, v, w, z, \theta, q, p) \in \mathcal{H} \left| \begin{array}{l} u \in H^2[0, L_0], v \in H^1[0, L_0], \\ w \in \tilde{V}^2[L_0, L], z \in \tilde{V}^1[L_0, L] \\ \theta \in V^2[0, L_0], q \in \tilde{V}^1[0, L_0] \\ p = v(L_0), u_x(0) = 0 \end{array} \right. \right\}.$$

Evolution Equation

System (2)–(5) can be rewritten as an evolutionary equation in \mathcal{H} :

$$\begin{cases} \frac{dU(t)}{dt} = \mathcal{A}U(t), & t > 0, \\ U(0) = U_0, \end{cases} \quad (6)$$

where

$$U(t) = (u, u_t, w, w_t, \theta, q, p)^T$$
$$U(0) = (u^0(x), u^1(x), w^0(x), w^1(x), \theta^0(x), q_0, p_0)^T \in \mathcal{H}.$$

Lemma

Let \mathcal{A} and \mathcal{H} be defined as in above. Then

- 1) \mathcal{A} is dissipative in \mathcal{H} ;
- 2) $0 \in \rho(\mathcal{A})$ and \mathcal{A}^{-1} is compact on \mathcal{H} . Therefore, the spectrum of \mathcal{A} consists of isolated eigenvalues of finite multiplicities only, i.e.,
 $\sigma(\mathcal{A}) = \sigma_p(\mathcal{A})$.

Theorem

Let \mathcal{A} and \mathcal{H} be defined as above. Then \mathcal{A} generates a C_0 semigroup of contractions $S(t)$ on \mathcal{H} .

Spectral Analysis of \mathcal{A}

For any $\lambda \in \sigma(\mathcal{A})$, suppose that $(u, v, w, z, \theta, q, p)^T \in \mathcal{D}(\mathcal{A})$ is an eigenvector of \mathcal{A} corresponding to λ .

$$\text{Eigenvalue Problem: } (\lambda I - \mathcal{A})(u, v, w, z, \theta, q, p)^T = 0$$



$$\left\{ \begin{array}{l} \lambda^2 u(x) - \alpha u_{xx}(x) + \beta \theta_x(x) = 0, x \in (0, L_0), \\ \lambda \theta(x) + r q_x(x) + \delta \lambda u_x(x) = 0, x \in (0, L_0), \\ (\tau \lambda + 1) q(x) + k \theta_x(x) = 0, x \in (0, L_0), \\ \lambda^2 w(x) - b w_{xx}(x) = 0, x \in (L_0, L), \\ u_x(0) = \theta(0) = w(L) = 0, \\ q(L_0) = 0, u(L_0) = w(L_0), \\ -M \lambda^2 u(L_0) = \alpha u_x(L_0) - \beta \theta(L_0) - b w_x(L_0), \quad M > 0. \end{array} \right. \quad (7)$$

Spectral analysis of \mathcal{A}

Lemma

Let \mathcal{A} and \mathcal{H} be defined as in above. Then there is no eigenvalue on the imaginary axis.

Lemma

If τ satisfies

$$\frac{M}{\tau} = \sqrt{\alpha} \frac{e^{\sqrt{\frac{1}{\alpha}} \frac{1}{\tau} L_0} - e^{-\sqrt{\frac{1}{\alpha}} \frac{1}{\tau} L_0}}{e^{\sqrt{\frac{1}{\alpha}} \frac{1}{\tau} L_0} + e^{-\sqrt{\frac{1}{\alpha}} \frac{1}{\tau} L_0}} - \sqrt{b} \frac{1 + e^{2\sqrt{\frac{1}{b}} \frac{1}{\tau} (L-L_0)}}{1 - e^{2\sqrt{\frac{1}{b}} \frac{1}{\tau} (L-L_0)}}, \quad (8)$$

then $-\frac{1}{\tau}$ is an eigenvalue of \mathcal{A} .

Spectral analysis of \mathcal{A}

Theorem

Let \mathcal{A} and \mathcal{H} be defined as in above. Then when $\lambda \neq -\frac{1}{\tau}$ with sufficiently large $|\lambda|$, the spectrum of \mathcal{A} has three branches given as follows

$$\lambda_1^n = \frac{1}{\sqrt{\xi_1}} \left(-\tilde{a}_1 + \frac{(n + \frac{1}{2})\pi i}{L_0} \right) + \mathcal{O}(n^{-2}), \quad n \in \mathbb{Z}, \quad (9)$$

$$\lambda_2^n = \frac{1}{\sqrt{\xi_2}} \left(\frac{\tilde{a}_2}{h} - \frac{(n + \frac{1}{2})\pi i}{L_0} \right) + \mathcal{O}(n^{-2}), \quad n \in \mathbb{Z}, \quad (10)$$

$$\lambda_3^n = \frac{n\sqrt{b}\pi i}{L - L_0} + \mathcal{O}(n^{-6}), \quad n \in \mathbb{Z}. \quad (11)$$

where $h < 0$ and $\tilde{a}_j > 0$, $j = 1, 2$.

spectral property

- there is no eigenvalue on imaginary axis;
- the imaginary axis is one asymptote of the spectrum;



Theorem

The total energy of the system (2)–(5) decays to 0 asymptotically.

Lemma

(Borichev and Tomilov [1] and Liu and Rao [2]) A C_0 semigroup $e^{t\mathcal{A}}$ of contractions on a Hilbert space satisfies

$$\|e^{t\mathcal{A}}U_0\| \leq Ct^{-\frac{1}{\delta}}\|U_0\|_{\mathcal{D}(\mathcal{A})}, \quad \forall U_0 \in \mathcal{D}(\mathcal{A}), \quad t \rightarrow \infty$$

for some constant $C > 0$, if the following conditions hold:

- 1). $\{i\beta | \beta \in \mathbb{R}\} \subset \rho(\mathcal{A})$
- 2). $\limsup_{|\beta| \rightarrow \infty} \frac{1}{|\beta|^\delta} \|(i\beta - \mathcal{A})^{-1}\| < \infty$.

Theorem

The $(T(t))_{t \geq 0}$ associated with the system (2)–(5) decays polynomially as

$$\|T(t)U_0\| \leq \frac{C}{t^{\frac{1}{\delta}}}\|U_0\|_{\mathcal{D}(\mathcal{A})}.$$

Question

Is $\frac{1}{6}$ a sharp estimate for the polynomial decay rate?

By one branch of the spectrum

$$\lambda_3^n = \frac{n\sqrt{b}\pi i}{L - L_0} + \mathcal{O}(n^{-6}), \quad n \in \mathbb{Z},$$

Polynomial decay order α must satisfy $\Re \lambda \geq \frac{C}{(\Re \lambda)^\alpha}$, $\lambda \in \sigma(A)$.

Theorem

$\frac{1}{6}$ is the optimal polynomial decay order.

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Welcome to Tianjin!

Thank you!