

# TWO-SIDED ESTIMATES OF FUNCTIONAL TYPE FOR GLOBAL AND LOCAL ERROR CONTROL

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## Error Estimate of Residual Type

These estimates have the following form:

$$\begin{aligned}
 & \|\nabla(u - u_h)\|_{0,\Omega} \leq \\
 & \leq C_0 \left( \left( \sum_{i=1}^N C_{1,i}^2 \text{diam}^2(\Delta_i) \|f + \text{div}\sigma_h\|_{0,\Delta_i}^2 \right)^{1/2} + \right. \\
 & \left. + \left( \sum_{k=1}^K \sum_{l>k, E_{kl} \subset \Omega} C_{2,kl}^2 |E_{kl}| \|j(\sigma_h \cdot \nu_{kl})\|_{0,E_{kl}}^2 \right)^{1/2} \right) \quad (\bullet)
 \end{aligned}$$

- In the above, the constant  $C_0$  strongly depend on properties of the concrete mesh  $\mathcal{T}_h$
- In fact, the most difficult type of error  $\varepsilon_3$  is neglected

**Quasi-Uniform Meshes:** In this case all generic constants  $C_{1,i}$  are approximately of the same value (and, similarly, the constants  $C_{2,kl}$ ), i.e., they can be replaced by only two constants  $c_1$  and  $c_2$ , respectively. Then RHS of the estimate (•) has a form

$$C \left( \sum_{i=1}^N \eta^2(\Delta_i) \right)^{1/2}$$

where

$$\eta^2(\Delta_i) = c_1^2 \text{diam}^2(\Delta_i) \|f + \text{div} \sigma_h\|_{0,\Delta_i}^2 + \frac{c_2^2}{2} \sum_{E_{kl} \subset \partial \Delta_i} |E_{kl}| \|j(\sigma_h \cdot \nu_{kl})\|_{0,E_{kl}}^2$$

and the constant  $C$  depends on  $c_1$ ,  $c_2$ , and  $C_0$

- However, for strongly nonhomogeneous meshes this estimate cannot be properly computed or, at least, lead to high overestimation of the error

**Arbitrary Meshes:** In the general case, sharper evaluation of the interpolation constants is desired. It leads to considerations of the following two variational problems

$$\inf_{v \in H_0^1(\Omega), v \neq 0} \frac{\|v\|_{1, \Omega_N(\Delta_i)}}{\|v - \pi_h v\|_{0, \Delta_i}} \text{diam}(\Delta_i)$$

and

$$\inf_{v \in H_0^1(\Omega), v \neq 0} \frac{\|v\|_{1, \Omega_N(E_{kl})}}{\|v - \pi_h v\|_{0, E_{kl}}} |E_{kl}|^{1/2}$$

- The constants  $C_{1,i}$  and  $C_{2,kl}$  are very difficult to compute or even to estimate from above, see, e.g., the paper [Carstensen, Funken]
- Moreover, the number of those constants depends on the dimension of the finite-element space  $V_h$  and can be very large
- *In fact, the estimation derived is only error indicator*

## Error Estimates Based on Post-Processing

These estimates have the following forms:

$$\|\nabla(u - u_h)\|_{0,\Omega} \approx \|G_h(\nabla u_h) - \nabla u_h\|_{0,\Omega}$$

$$\ell(u - u_h) \approx E_0(u_h, v_\tau) + \int_{\Omega} (G_h(\nabla u_h) - \nabla u_h) \cdot (G_\tau(\nabla v_\tau) - \nabla v_\tau) dx$$

where  $G_h$  and  $G_\tau$  are suitable gradient averaging procedures

- The above error indicators are quite efficient and reliable in many situations, but there are problems (even simple ones) for which they completely fail

**Test 1:** In this test we demonstrate that the global error indicator completely fails if the problem data is not sufficiently smooth

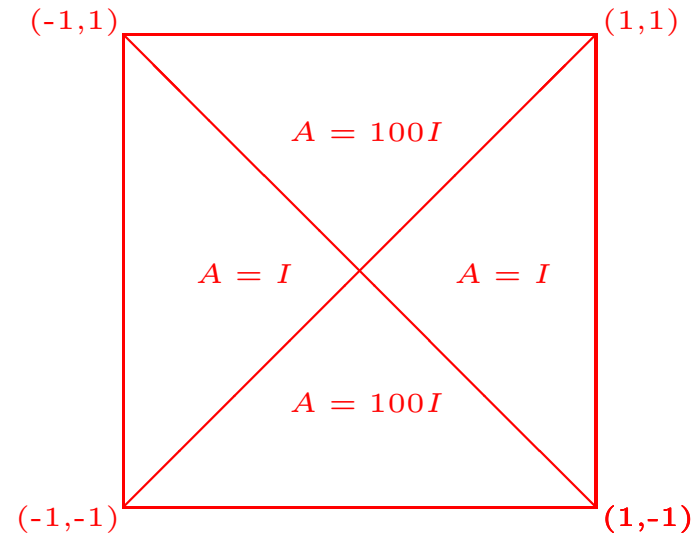


Figure 1: Definition of the coefficient matrix  $\mathcal{A}$

For the indicator value we have 17.9628, which is quite different from the error  $\|\nabla(u - u_h)\|_{\Omega}^2 \approx 0.8584$

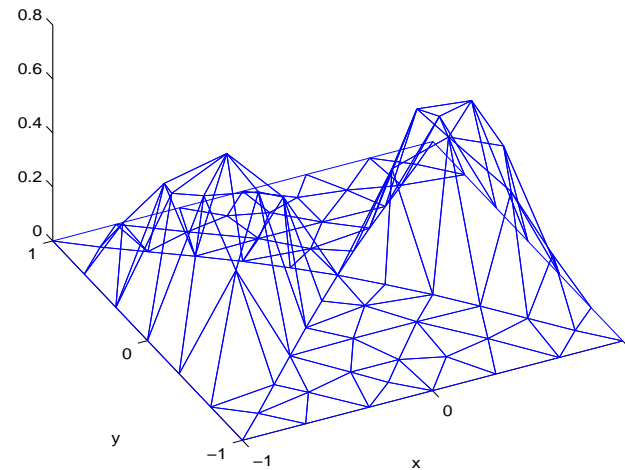
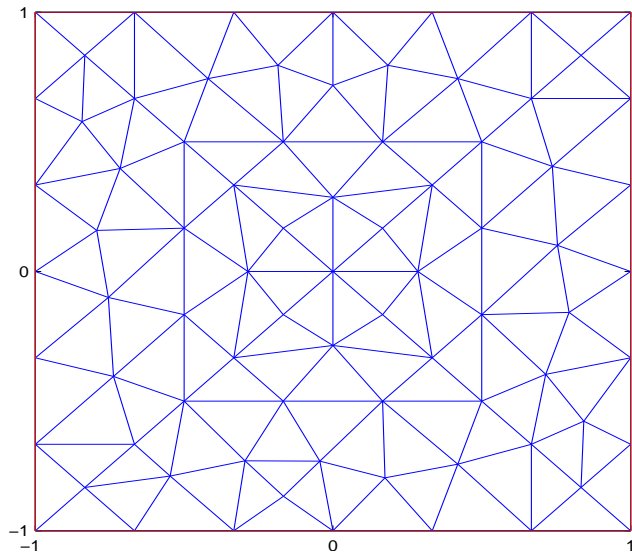
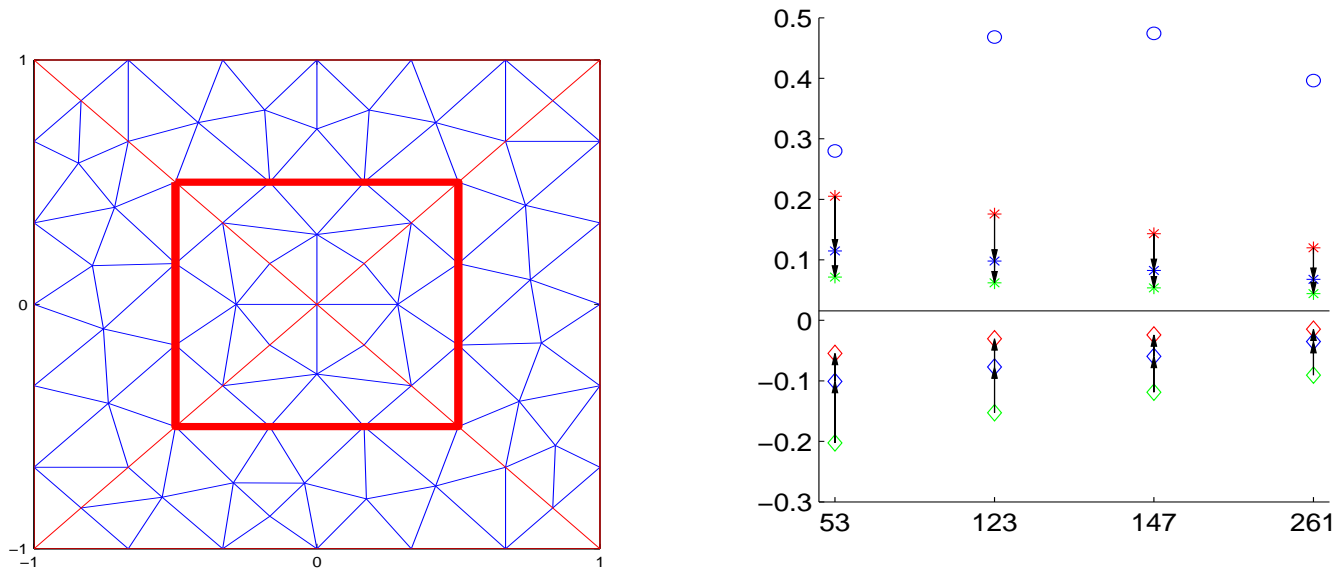


Figure 2: Solution domain  $\Omega$  with computational mesh  $\mathcal{T}_h$  having 78 nodes, the finite element approximation  $u_h$

**Test 2:** To demonstrate that the local indicator fails, we take the same problem as in **Test 1**



We performed error estimation using several different meshes for the adjoint problem. We clearly see that the values of indicator (by circles) are essentially bigger than  $\ell(u - u_h) = 0.0157$  for all choices of the mesh for the adjoint problem



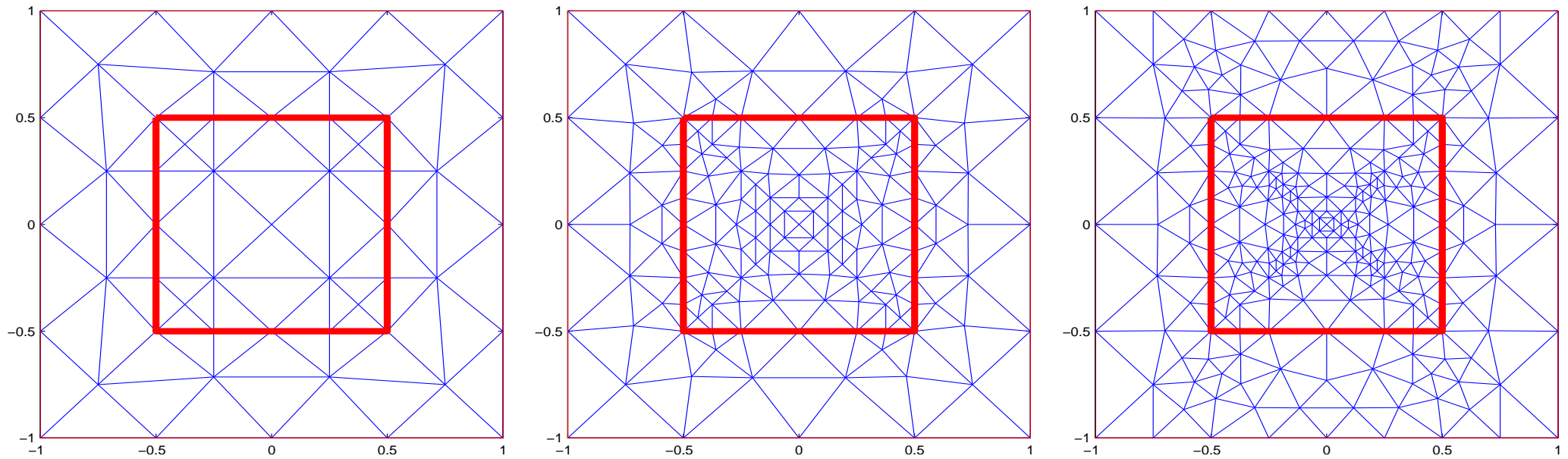


Figure 3: Meshes for adjoint problem with 53, 147, and 261 nodes

**NEW TECHNOLOGIES FOR ERROR**

**CONTROL DURING COMPUTATIONS:**

**THEORETICAL PART**

## Main Targets of New Technologies

To create a set of numerical algorithms (and respective computer codes), which is able to give the following information about the accuracy of obtained approximate solutions:

- guaranteed two-sided bounds of the error in the global norm
- guaranteed two-sided estimation of various local errors
- guidelines for efficient mesh adaptation in order to decrease the errors
- All error estimates are required to be absolutely independent of the nature of the method used to obtain approximations, since we never have in practice e.g. true finite element solutions

## Model Elliptic Type BVPs

Find a function  $u$  such that

$$-\operatorname{div}(A\nabla u) + cu = f \quad \text{in } \Omega \quad (1)$$

$$u = u_0 \quad \text{on } \Gamma_D \quad (2)$$

$$\nu^T \cdot A\nabla u = g \quad \text{on } \Gamma_N \quad (3)$$

where  $\Omega$  is bounded domain in  $\mathbf{R}^d$  with Lipschitz boundary  
 $\partial\Omega = \bar{\Gamma}_D \cup \bar{\Gamma}_N$ ,  $\operatorname{meas}_{d-1}\Gamma_D > 0$ ,  $\nu$  is the outward normal to  $\partial\Omega$

**Assumptions:**  $f \in L_2(\Omega)$ ,  $u_0 \in H^1(\Omega)$ ,  $g \in L_2(\Gamma_N)$ ,  $c \in L_\infty(\Omega)$ , matrix of coefficients  $A$  is symmetric, with entries  $a_{ij} \in L_\infty(\Omega)$ ,  $i, j = 1, \dots, d$ , and is such that

$$C_2|\xi|^2 \geq A(x)\xi \cdot \xi \geq C_1|\xi|^2 \quad \forall \xi \in \mathbf{R}^d \quad \forall x \in \Omega \quad (4)$$

In addition, let us assume that almost everywhere in  $\Omega$

$$c \geq 0 \quad (5)$$

and stand the denotation

$$\Omega^c := \text{supp } c = \{x \in \Omega \mid c(x) > 0\} \quad (6)$$

**Weak formulation:** Find  $u \in u_0 + H_{\Gamma_D}^1(\Omega)$  such that

$$\int_{\Omega} A \nabla u \cdot \nabla w \, dx + \int_{\Omega} c u w \, dx = \int_{\Omega} f w \, dx + \int_{\Gamma_N} g w \, ds \quad \forall w \in H_{\Gamma_D}^1(\Omega) \quad (7)$$

where

$$H_{\Gamma_D}^1(\Omega) := \{v \in H^1(\Omega) \mid v = 0 \text{ on } \Gamma_D\} \quad (8)$$

- Existence and uniqueness of the *weak solution* defined by (7) is provided by the well-known Lax-Milgram lemma

We can define *bilinear and linear forms* as follows

$$a(v, w) := \int_{\Omega} A \nabla v \cdot \nabla w \, dx + \int_{\Omega} cvw \, dx, \quad F(w) := \int_{\Omega} fw \, dx + \int_{\Gamma_N} gw \, ds \quad (9)$$

- Then formulation (7) can be written in a more compact form:

$$\text{Find } u \in u_0 + H_{\Gamma_D}^1(\Omega) \text{ such that } a(u, w) = F(w) \quad \forall w \in H_{\Gamma_D}^1(\Omega) \quad (\dagger)$$

- Lax-Milgram lemma is, in fact, applied to such a problem:

$$\text{Find } u_{\star} \in H_{\Gamma_D}^1(\Omega) \text{ such that } a(u_{\star}, w) = \bar{F}(w) \quad \forall w \in H_{\Gamma_D}^1(\Omega) \quad (\star)$$

where  $\bar{F}(w) = F(w) - a(u_0, w)$ . And then we take  $u = u_0 + u_{\star}$

- Formulations  $(\dagger)$  and  $(\star)$  are equivalent

## Energy Functional

Having formulation  $(\star)$ , it is possible to define a convenient *energy functional*  $J$  as follows

$$J(w) := \frac{1}{2}a(w, w) - \bar{F}(w) \quad (10)$$

- The term “*energy functional*” comes from the fact that the functional taken in the form (10) often presents the total energy of the system under analysis
- It is well-known that problem  $(\star)$  is equivalent to the problem of finding the minimizer of the energy functional over space  $H_{\Gamma_D}^1(\Omega)$



## Control of Computational Errors

- Let  $\bar{u} = u_0 + \bar{u}_*$ , where  $\bar{u}_* \in H_{\Gamma_D}^1(\Omega)$ , be *any function* from the set  $u_0 + H_{\Gamma_D}^1(\Omega)$  (e.g., computed by some numerical method) considered as approximation of  $u$
- We are interested in *diverse and reliable estimation and control of the deviation (or error)  $e := u - \bar{u} = u_* - \bar{u}_*$*  during computations
- So far *two ways* of the error control *aimed at two different final goals of the whole computational process* are well studied and widely used in mathematical and engineering communities
- Roughly speaking, those are – *control of the overall accuracy* of obtained approximations and *control of various local errors*

## Overall Error in Energy Norm

First of all, we can easily prove that

$$a(u - \bar{u}, u - \bar{u}) = a(u_\star - \bar{u}_\star, u_\star - \bar{u}_\star) = 2(J(\bar{u}_\star) - J(u_\star)) \quad (11)$$

which means that the value

$$a(u - \bar{u}, u - \bar{u}) = \int_{\Omega} A \nabla(u - \bar{u}) \cdot \nabla(u - \bar{u}) dx + \int_{\Omega} c(u - \bar{u})^2 dx \quad (12)$$

is a natural measure for the overall quality of approximation  $\bar{u}$

- Since  $\sqrt{a(\cdot, \cdot)}$  is called the *energy norm*, the estimation of value (12) is often called the *error estimation in the energy norm*

- In the case  $c \equiv 0$ , the value (12) reduces to the well-known form

$$\int_{\Omega} A \nabla(u - \bar{u}) \cdot \nabla(u - \bar{u}) dx \quad (13)$$

which is nowadays widely used as a measure of accuracy of  $\bar{u}$

- The error control in the global energy norm (13) has been in a focus of active research work since the works by Babuška and Rheinboldt in 1978

## Control of Local Errors

- Another trend in a posteriori error estimation is based on concept of local error control in addition to classical control in energy norm
- This approach is motivated by practical needs: analysts are often interested not only in the value of the overall error, but also in *controlling errors over certain parts of solution domain*, or relative to some interesting characteristics (*“quantities of interest”*)
- Common way for performing such type control is to introduce *linear functional  $\ell$*  associated with subdomain of interest (and/or with a “quantity of interest”) and to obtain a computable estimate for  $\ell(u - \bar{u})$

- For example, one can be interested in estimation of

$$\ell(e) = \ell(u - \bar{u}) = \int_{\Omega} \varphi(u - \bar{u}) \, dx \quad (14)$$

where  $\varphi \in L_2(\Omega)$  and  $\text{supp } \varphi = \omega \subseteq \Omega$ , which provides us with info on error  $e$  locally, in subdomain  $\omega$

- Estimates for  $\ell(u - \bar{u})$  can also be used for estimation of unknown “quantity of interest”  $\ell(u)$ , since

$$\ell(u) = \ell(\bar{u}) + \ell(u - \bar{u})$$

where  $\ell(\bar{u})$  is computable and  $\ell(u - \bar{u})$  is estimated

- Certain “quantities of interest” (e.g., in linear elasticity) are often more interesting for the practitioners than the solution itself

## Inequalities and Constants

- In what follows we need Friedrichs' inequality

$$\|w\|_{0,\Omega} \leq C_{\Omega,\Gamma_D} \|\nabla w\|_{0,\Omega} \quad \forall w \in H_{\Gamma_D}^1(\Omega) \quad (15)$$

and the inequality in the trace theorem

$$\|w\|_{0,\partial\Omega} \leq C_{\partial\Omega} \|w\|_{1,\Omega} \quad \forall w \in H^1(\Omega) \quad (16)$$

where  $C_{\Omega,\Gamma_D}$  and  $C_{\partial\Omega}$  are positive constants depending only on  $\Omega$ ,  $\Gamma_D$ , and  $\partial\Omega$

- $\|\cdot\|_{0,\Omega}$  and  $\|\cdot\|_{1,\Omega}$  stand for norms in  $L_2(\Omega)$  and  $H^1(\Omega)$
- $\|\cdot\|_{0,\partial\Omega}$  is the norm in  $L_2(\partial\Omega)$

## Two-Sided Estimates of Error in Energy Norm

- Let  $\chi_S$  be the characteristic function of set  $S$ :  $\chi_S(x) := 1$  if  $x \in S$ , and  $\chi_S(x) := 0$  if  $x \notin S$
- Let  $\|y\|_\Omega := \sqrt{\int_\Omega Ay \cdot y \, dx}$  for  $y \in L_2(\Omega, \mathbf{R}^d)$
- We shall need two above defined constants  $C_{\Omega, \Gamma_D}$ ,  $C_{\partial\Omega}$ , and also

$$C_{\Omega, \partial\Omega} := \frac{C_{\partial\Omega} \sqrt{1 + C_{\Omega, \Gamma_D}^2}}{\sqrt{C_1}}$$

- $H_N(\Omega, \text{div}) := \{y \in L_2(\Omega, \mathbf{R}^d) \mid \text{div } y \in L_2(\Omega), \nu^T \cdot y \in L_2(\Gamma_N)\}$
- $\Omega^c := \text{supp } c = \{x \in \Omega \mid c(x) > 0\}$

**Upper estimate:** We have the following upper estimate

$$\begin{aligned}
a(u - \bar{u}, u - \bar{u}) &\leq \left\| \frac{1}{\sqrt{c}}(f + \operatorname{div} y^* - c\bar{u}) \right\|_{0, \Omega^c}^2 + \\
+(1 + \alpha) &\|A^{-1}y^* - \nabla \bar{u}\|_{\Omega}^2 + (1 + \frac{1}{\alpha})(1 + \beta) \frac{C_{\Omega, \Gamma_D}^2}{C_1} \|f + \operatorname{div} y^*\|_{0, \Omega \setminus \bar{\Omega}^c}^2 \\
&+(1 + \frac{1}{\alpha})(1 + \frac{1}{\beta}) C_{\Omega, \partial\Omega}^2 \|g - \nu^T \cdot y^*\|_{0, \Gamma_N}^2
\end{aligned} \tag{17}$$

where  $\alpha, \beta$  are any positive numbers and  $y^*$  is any function from  $H_N(\Omega, \operatorname{div})$



**P r o o f :** First of all, we notice that

$$a(u - \bar{u}, u - \bar{u}) = \|\nabla(u - \bar{u})\|_{\Omega}^2 + \|\sqrt{c}(u - \bar{u})\|_{0, \Omega^c}^2 \quad (18)$$

Further, using fact that  $u - \bar{u} \in H_{\Gamma_D}^1(\Omega)$ , integral identity (7), Green formula, and simple regrouping of terms in below we get

$$\begin{aligned} a(u - \bar{u}, u - \bar{u}) &= \int_{\Omega} f(u - \bar{u}) dx + \int_{\Gamma_N} g(u - \bar{u}) ds - \int_{\Omega} A \nabla \bar{u} \cdot \nabla (u - \bar{u}) dx \\ &\quad - \int_{\Omega} c \bar{u} (u - \bar{u}) dx = \int_{\Omega} (f - c \bar{u})(u - \bar{u}) dx + \int_{\Gamma_N} g(u - \bar{u}) ds \\ &\quad - \int_{\Omega} (A \nabla \bar{u} - y^*) \cdot \nabla (u - \bar{u}) dx - \int_{\Omega} y^* \cdot \nabla (u - \bar{u}) dx = \end{aligned} \quad (19)$$

$$\begin{aligned}
&= \int_{\Omega} (f + \operatorname{div} y^* - c\bar{u})(u - \bar{u}) \, dx - \int_{\Omega} A(\nabla\bar{u} - A^{-1}y^*) \cdot \nabla(u - \bar{u}) \, dx \\
&\quad + \int_{\Gamma_N} g(u - \bar{u}) \, ds - \int_{\Gamma_N} \nu^T \cdot y^*(u - \bar{u}) \, ds = \\
&= \int_{\Omega} A(A^{-1}y^* - \nabla\bar{u}) \cdot \nabla(u - \bar{u}) \, dx + \int_{\Omega} (f + \operatorname{div} y^* - c\bar{u})(u - \bar{u}) \, dx \\
&\quad + \int_{\Gamma_N} (g - \nu^T \cdot y^*)(u - \bar{u}) \, ds
\end{aligned}$$

where  $y^*$  is any function from the space  $H_N(\Omega, \operatorname{div})$

Now, the right-hand side (RHS) of the above equality can be estimated, using the Cauchy-Schwarz inequality, denotation  $\Omega^c$  and trace theorem from above as follows

$$\begin{aligned}
\text{RHS (19)} &\leq \|A^{-1}y^* - \nabla\bar{u}\|_{\Omega} \|\nabla(u - \bar{u})\|_{\Omega} + \|g - \nu^T \cdot y^*\|_{0, \Gamma_N} \|u - \bar{u}\|_{0, \Gamma_N} \\
&\quad + \int_{\Omega} (f + \operatorname{div} y^* - c\bar{u})(u - \bar{u}) \, dx \leq \\
&\leq \|A^{-1}y^* - \nabla\bar{u}\|_{\Omega} \|\nabla(u - \bar{u})\|_{\Omega} + \|g - \nu^T \cdot y^*\|_{0, \Gamma_N} C_{\partial\Omega} \|u - \bar{u}\|_{1, \Omega} \quad (20) \\
&+ \int_{\Omega^c} \frac{1}{\sqrt{c}} (f + \operatorname{div} y^* - c\bar{u}) \sqrt{c}(u - \bar{u}) \, dx + \int_{\Omega \setminus \bar{\Omega}^c} (f + \operatorname{div} y^* - c\bar{u})(u - \bar{u}) \, dx
\end{aligned}$$

Further, using the ellipticity condition, Friedrichs' inequality, and the inequality

$$|ab| \leq \frac{1}{2}a^2 + \frac{1}{2}b^2 \quad (21)$$

we observe that

$$\begin{aligned} & \text{RHS (20)} \leq \\ & \leq \left( \|A^{-1}y^* - \nabla \bar{u}\|_{\Omega} + \frac{C_{\partial\Omega} \sqrt{1 + C_{\Omega, \Gamma_D}^2}}{\sqrt{C_1}} \|g - \nu^T \cdot y^*\|_{0, \Gamma_N} \right) \|\nabla(u - \bar{u})\|_{\Omega} \\ & \quad + \frac{1}{2} \|\sqrt{c}(u - \bar{u})\|_{0, \Omega^c}^2 + \frac{1}{2} \left\| \frac{1}{\sqrt{c}} (f + \text{div } y^* - c\bar{u}) \right\|_{0, \Omega^c}^2 \\ & \quad + \int_{\Omega} \chi_{\Omega \setminus \bar{\Omega}^c} (f + \text{div } y^* - c\bar{u}) (u - \bar{u}) \, dx \leq \end{aligned}$$

$$\begin{aligned}
&\leq \left( \|A^{-1}y^* - \nabla\bar{u}\|_{\Omega} + C_{\Omega,\partial\Omega} \|g - \nu^T \cdot y^*\|_{0,\Gamma_N} \right) \|\nabla(u - \bar{u})\|_{\Omega} \quad (22) \\
&\quad + \frac{1}{2} \|\sqrt{c}(u - \bar{u})\|_{0,\Omega^c}^2 + \frac{1}{2} \left\| \frac{1}{\sqrt{c}}(f + \operatorname{div} y^* - c\bar{u}) \right\|_{0,\Omega^c}^2 \\
&\quad + \|\chi_{\Omega \setminus \bar{\Omega}^c}(f + \operatorname{div} y^* - c\bar{u})\|_{0,\Omega} \|u - \bar{u}\|_{0,\Omega}
\end{aligned}$$

where  $C_{\Omega,\partial\Omega} := \frac{C_{\partial\Omega} \sqrt{1+C_{\Omega,\Gamma_D}^2}}{\sqrt{C_1}}$

Regrouping terms in RHS of (22) and using again the inequality (21), we get an estimate

$$\text{RHS (22)} \leq$$

$$\left( \left\| A^{-1}y^* - \nabla \bar{u} \right\|_{\Omega} + C_{\Omega, \partial\Omega} \left\| g - \nu^T \cdot y^* \right\|_{0, \Gamma_N} + \frac{C_{\Omega, \Gamma_D}}{\sqrt{C_1}} \left\| f + \text{div } y^* \right\|_{0, \Omega \setminus \bar{\Omega}^c} \right) \times$$

$$\times \left\| \nabla(u - \bar{u}) \right\|_{\Omega} + \frac{1}{2} \left\| \sqrt{c}(u - \bar{u}) \right\|_{0, \Omega^c}^2 + \frac{1}{2} \left\| \frac{1}{\sqrt{c}}(f + \text{div } y^* - c\bar{u}) \right\|_{0, \Omega^c}^2 \leq$$

$$\frac{1}{2} \left( \left\| A^{-1}y^* - \nabla \bar{u} \right\|_{\Omega} + C_{\Omega, \partial\Omega} \left\| g - \nu^T \cdot y^* \right\|_{0, \Gamma_N} + \frac{C_{\Omega, \Gamma_D}}{\sqrt{C_1}} \left\| f + \text{div } y^* \right\|_{0, \Omega \setminus \bar{\Omega}^c} \right)^2$$

$$(23)$$

$$+ \frac{1}{2} \left\| \nabla(u - \bar{u}) \right\|_{\Omega}^2 + \frac{1}{2} \left\| \sqrt{c}(u - \bar{u}) \right\|_{0, \Omega^c}^2 + \frac{1}{2} \left\| \frac{1}{\sqrt{c}}(f + \text{div } y^* - c\bar{u}) \right\|_{0, \Omega^c}^2$$

Using now (18) and the final inequality resulting from (19)–(20) and (22)–(23), multiplying it by two and regrouping, we immediately get for the error in the energy norm that

$$\begin{aligned}
& a(u - \bar{u}, u - \bar{u}) = \\
& = \|\nabla(u - \bar{u})\|_{\Omega}^2 + \|\sqrt{c}(u - \bar{u})\|_{0,\Omega^c}^2 \leq \left\| \frac{1}{\sqrt{c}}(f + \operatorname{div} y^* - c\bar{u}) \right\|_{0,\Omega^c}^2 \\
& + \left( \left\| A^{-1}y^* - \nabla\bar{u} \right\|_{\Omega} + \frac{C_{\Omega,\Gamma_D}}{\sqrt{C_1}} \|f + \operatorname{div} y^*\|_{0,\Omega \setminus \bar{\Omega}^c} + C_{\Omega,\partial\Omega} \|g - \nu^T \cdot y^*\|_{0,\Gamma_N} \right)^2
\end{aligned} \tag{24}$$

Finally, using two times the inequality

$$(a + b)^2 \leq (1 + \lambda)a^2 + \left(1 + \frac{1}{\lambda}\right)b^2 \quad (\lambda > 0)$$

for terms in round brackets in (24), we get our upper estimate

**Lower estimate:** We have the following lower bound

$$a(u - \bar{u}, u - \bar{u}) \geq 2(J(\bar{u}_\star) - J(w)) \quad (25)$$

where  $w$  is any function from  $H_{\Gamma_D}^1(\Omega)$  and  $J$  is energy functional

**P r o o f :** First, we prove that

$$a(u - \bar{u}, u - \bar{u}) = a(u_\star - \bar{u}_\star, u_\star - \bar{u}_\star) = 2(J(\bar{u}_\star) - J(u_\star)) \quad (26)$$

Indeed, we observe

$$\begin{aligned} 2(J(\bar{u}_\star) - J(u_\star)) &= a(\bar{u}_\star, \bar{u}_\star) - 2\bar{F}(\bar{u}_\star) - a(u_\star, u_\star) + 2\bar{F}(u_\star) \\ &= a(\bar{u}_\star, \bar{u}_\star) - a(u_\star, u_\star) + 2\bar{F}(u_\star - \bar{u}_\star) \\ &= a(\bar{u}_\star, \bar{u}_\star) - a(u_\star, u_\star) + 2a(u_\star, u_\star - \bar{u}_\star) \\ &= a(\bar{u}_\star, \bar{u}_\star) + a(u_\star, u_\star) - 2a(u_\star, \bar{u}_\star) = a(u_\star - \bar{u}_\star, u_\star - \bar{u}_\star) \end{aligned}$$

Since  $u_\star$  minimizes  $J$ , we have  $J(u_\star) \leq J(w) \quad \forall w \in H_{\Gamma_D}^1(\Omega)$ , which proves (25)



## Comments on Estimates

- In order to derive the upper and lower estimates, we did not require the function  $\bar{u}$  be a finite element approximation (or computed by some another specific numerical method); it is simply any function from  $u_0 + H_{\Gamma_D}^1(\Omega)$
- The upper estimate cannot be improved. Indeed, if one takes  $y^* = A\nabla u$ , which obviously belongs to  $H_N(\Omega, \text{div})$ , then the last two terms in the right-hand side of (17) vanish. Further, taking  $\alpha \rightarrow 0$ , we finally observe that the inequality (17) holds as equality. To prove that the lower estimate cannot be improved either, we should, obviously, take  $w = u_\star \in H_{\Gamma_D}^1(\Omega)$

- The upper estimate contains only two global constants,  $C_{\Omega, \Gamma_D}$  and  $C_{\partial\Omega}$ , which do not depend on the computational process at all. They have to be computed (or accurately estimated from above) only once when the problem is posed
- In many works, devoted to a posteriori error estimation, one usually takes  $c \equiv 0$ . In this case  $a(u - \bar{u}, u - \bar{u}) = \|\nabla(u - \bar{u})\|_{\Omega}^2$ , the set  $\Omega^c = \emptyset$ , and the upper estimate takes a simpler form

$$\begin{aligned}
& a(u - \bar{u}, u - \bar{u}) \leq \\
& \leq (1 + \alpha) \|A^{-1}y^* - \nabla\bar{u}\|_{\Omega}^2 + (1 + \frac{1}{\alpha})(1 + \beta) \frac{C_{\Omega, \Gamma_D}^2}{C_1} \|f + \operatorname{div} y^*\|_{0, \Omega}^2 \\
& \quad + (1 + \frac{1}{\alpha})(1 + \frac{1}{\beta}) C_{\Omega, \partial\Omega}^2 \|g - \nu^T \cdot y^*\|_{0, \Gamma_N}^2 \quad (27)
\end{aligned}$$

- For the pure Dirichlet boundary condition, the third term in RHS of (27) does not exist, and, since the estimate is valid for any positive  $\beta$ , we can take it be zero. Then, we get

$$\begin{aligned}
 & a(u - \bar{u}, u - \bar{u}) \leq \\
 & \leq (1 + \alpha) \|A^{-1}y^* - \nabla \bar{u}\|_{\Omega}^2 + (1 + \frac{1}{\alpha}) \frac{C_{\Omega, \Gamma_D}^2}{C_1} \|f + \operatorname{div} y^*\|_{0, \Omega}^2 \quad (28)
 \end{aligned}$$

- The upper estimate (28) was first obtained in [Repin, 1997] using quite complicated tools of the duality theory. Later it was obtained in [Repin, Sauter, Smolianski, 2003] for simple Poisson equation, using an idea of the Helmholtz decomposition of  $L_2(\Omega, \mathbf{R}^d)$ . The estimate (27) is derived in [Repin, Sauter, Smolianski, 2004] using again the duality theory

- Our approach for derivation of the estimates is different from those used in the above mentioned works and is simpler
- More general problems with mixed Dirichlet/Neumann/Robin (D/N/R) boundary conditions are treated in a just presented way in [Korotov, 2006], where (17) follows as a particular case
- In the case of pure Dirichlet (D) condition or Dirichlet/Robin (D/R) mixed boundary condition we have to compute, or estimate from above, only one constant  $C_{\Omega, \Gamma_D}$
- If, in addition to D or D/R boundary conditions, we have  $c(x) \geq c_0 > 0$  in  $\Omega$  we do not need any global constants for our estimation at all

## Two-Sided Estimates for Local Errors

- The second technology is developed for controlling local errors
- To demonstrate the main idea of the approach and avoid unnecessary technical details we shall consider only the simplest problem. More general cases can be treated just in the same fashion

**Model Problem:** Find a function  $u$  such that

$$-\operatorname{div}(A\nabla u) = f \quad \text{in } \Omega \quad (29)$$

$$u = 0 \quad \text{on } \partial\Omega \quad (30)$$

- I.e.,  $c \equiv 0$ ,  $u_0 = 0$ ,  $\Gamma_N = \emptyset$ , and  $a(u - \bar{u}, u - \bar{u}) = \|\nabla e\|_{\Omega}^2$  then

**Weak formulation:** Find  $u \in H_0^1(\Omega)$  such that

$$\int_{\Omega} A \nabla u \cdot \nabla w \, dx = \int_{\Omega} f w \, dx \quad \forall w \in H_0^1(\Omega) \quad (31)$$

- We want to estimate the deviation  $e = u - \bar{u}$  measured in terms of suitable linear functional  $\ell$

Example: 
$$\ell(e) = \int_{\Omega} \varphi (u - \bar{u}) \, dx \quad (32)$$

where  $\varphi \in L_2(\Omega)$  and  $\text{supp } \varphi = \omega \subseteq \Omega$

**Adjoint problem (AP):** Find  $v \in H_0^1(\Omega)$  such that

$$\int_{\Omega} A \nabla v \cdot \nabla w \, dx = \ell(w) \quad \forall w \in H_0^1(\Omega) \quad (33)$$

- Adjoint problem is uniquely solvable

- Usually we only have some approximation  $\bar{v}$  (e.g., computed by FEM on *adjoint mesh*  $\mathcal{T}_\tau$ )
- It can be shown that

$$\ell(u - \bar{u}) = E_0(\bar{u}, \bar{v}) + E_1(e, e_\ell) \quad (34)$$

where

$$E_0(\bar{u}, \bar{v}) = \int_{\Omega} f \bar{v} \, dx - \int_{\Omega} A \nabla \bar{v} \cdot \nabla \bar{u} \, dx \quad (35)$$

and

$$E_1(e, e_\ell) = \int_{\Omega} A \nabla e \cdot \nabla e_\ell \, dx \quad (36)$$

- Denotation  $e_\ell := v - \bar{v}$  stands for computational error in **AP**

Proof: From the integral equality for **AP**, we have ( $w := u - \bar{u}$ )

$$\begin{aligned} \ell(u - \bar{u}) &= \int_{\Omega} A \nabla v \cdot \nabla(u - \bar{u}) \, dx = \\ &= \int_{\Omega} A \nabla(v - \bar{v}) \cdot \nabla(u - \bar{u}) \, dx + \int_{\Omega} A \nabla \bar{v} \cdot \nabla(u - \bar{u}) \, dx = \\ &= E_1(e, e_\ell) - \int_{\Omega} A \nabla \bar{v} \cdot \nabla \bar{u} \, dx + \int_{\Omega} A \nabla u \cdot \nabla \bar{v} \, dx = \\ &= E_1(e, e_\ell) - \int_{\Omega} A \nabla \bar{v} \cdot \nabla \bar{u} \, dx + \int_{\Omega} f \bar{v} \, dx = \\ &= E_0(\bar{u}, \bar{v}) + E_1(e, e_\ell) \end{aligned}$$



- First term  $E_0$  is directly computable once  $\bar{u}$  and  $\bar{v}$  are computed, but term  $E_1$  contains unknown gradients  $\nabla u$  and  $\nabla v$
- In order to estimate  $E_1$ , we use the following identity

$$2E_1(e, e_\ell) = \|\nabla(\gamma e + \frac{1}{\gamma}e_\ell)\|_\Omega^2 - \gamma^2 \|\nabla e\|_\Omega^2 - \frac{1}{\gamma^2} \|\nabla e_\ell\|_\Omega^2 \quad (37)$$

valid for any  $\gamma > 0$

- The above identity contains errors in energy norm for both, primal and adjoint problems, and we can use the corresponding two-sided estimates

$$M_f^\ominus \leq \|\nabla e\|_\Omega^2 \leq M_f^\oplus \quad \& \quad M_\ell^\ominus \leq \|\nabla e_\ell\|_\Omega^2 \leq M_\ell^\oplus \quad (38)$$

- For the first term in RHS of (37), we observe that

$$\|\nabla(\gamma e + \frac{1}{\gamma} e_\ell)\|_\Omega^2 = \|\nabla(\gamma u + \frac{1}{\gamma} v) - \nabla(\gamma \bar{u} + \frac{1}{\gamma} \bar{v})\|_\Omega^2 \quad (39)$$

- Function  $\gamma u + \frac{1}{\gamma} v$  can be perceived as the solution of problem:

Find  $u_\gamma \in H_0^1(\Omega)$  such that

$$\int_\Omega A \nabla u_\gamma \cdot \nabla w \, dx = \gamma \int_\Omega f w \, dx + \frac{1}{\gamma} \ell(w) \quad \forall w \in H_0^1(\Omega) \quad (40)$$

and function  $\gamma \bar{u} + \frac{1}{\gamma} \bar{v}$  can be considered as an approximation of  $u_\gamma$

- Then we can again use two-sided estimates

$$M_{\gamma,f,\ell}^\ominus \leq \|\nabla(\gamma e + \frac{1}{\gamma} e_\ell)\|_\Omega^2 \leq M_{\gamma,f,\ell}^\oplus \quad (41)$$

- We immediately observe that

$$\frac{1}{2}(M_{\gamma,f,\ell}^{\ominus} - \gamma^2 M_f^{\oplus} - \frac{1}{\gamma^2} M_{\ell}^{\oplus}) \leq E_1(e, e_{\ell})$$

and

$$E_1(e, e_{\ell}) \leq \frac{1}{2}(M_{\gamma,f,\ell}^{\oplus} - \gamma^2 M_f^{\ominus} - \frac{1}{\gamma^2} M_{\ell}^{\ominus})$$

which together with directly computable term

$$E_0(\bar{u}, \bar{v}) = \int_{\Omega} f \bar{v} \, dx - \int_{\Omega} \mathbf{A} \nabla \bar{v} \cdot \nabla \bar{u} \, dx$$

provide with two-sided estimates for the error  $\ell(e) = E_0 + E_1$

## Mesh Adaptivity

- For both types of error control, the proposed upper and lower estimates have integral form, i.e, each one can be presented as a single integral over the solution domain  $\Omega$
- Such a form suggests a standard mesh adaptivity procedure based on refining those elements of the mesh which give a large contribution to the total error

**NEW TECHNOLOGIES:**

**PRACTICAL REALIZATION**

**AND NUMERICAL TESTS**

## On Computation of Global Constants

- Our estimates contains only two unknown constants  $C_{\Omega, \Gamma_D}, C_{\partial\Omega}$
- Estimation of those two constants from above is sufficient
- Other estimation techniques involve *many unknown constants*:
  - which are very hard to estimate
  - which have to be always recomputed if the mesh changes
- But  $C_{\Omega, \Gamma_D}, C_{\partial\Omega}$  remain the same under any changes of meshes

## Series of Meshes

- We often perform computations on a series of meshes  $\mathcal{T}_{h_1}, \mathcal{T}_{h_2}, \mathcal{T}_{h_3}, \dots$ , where  $h = h_1 > h_2 > h_3 > \dots$
- Thus, we always have several successive approximations  $u_{h_1}, u_{h_2}, u_{h_3}, \dots$ , where  $u_{h_i} = u_0 + \bar{u}_{h_i}$
- We normally try to have  $J(\bar{u}_{h_1}) > J(\bar{u}_{h_2}) > J(\bar{u}_{h_3}) > \dots$

## On Computation of Upper Bounds

- For simplicity we consider the minimization issue under condition that approximation is computed by FEM, i.e.,  $\bar{u} \equiv u_h$

$$M^\oplus = M^\oplus(u_h, \alpha, \beta, y^*) \quad (42)$$

- Coarse upper bounds can be computed fast using, e.g.,  $y^* = G_\mu(\nabla u_\mu) \in H_N(\Omega, \text{div})$ ,  $\mu = h_1, h_2, h_3, \dots$ , and  $G_\mu$  is some gradient averaging operator
- Sharp estimates require a real minimization of  $M^\oplus(u_h, \alpha, \beta, y^*)$  with respect to “free variables”  $y^*$  and  $\alpha$  &  $\beta$



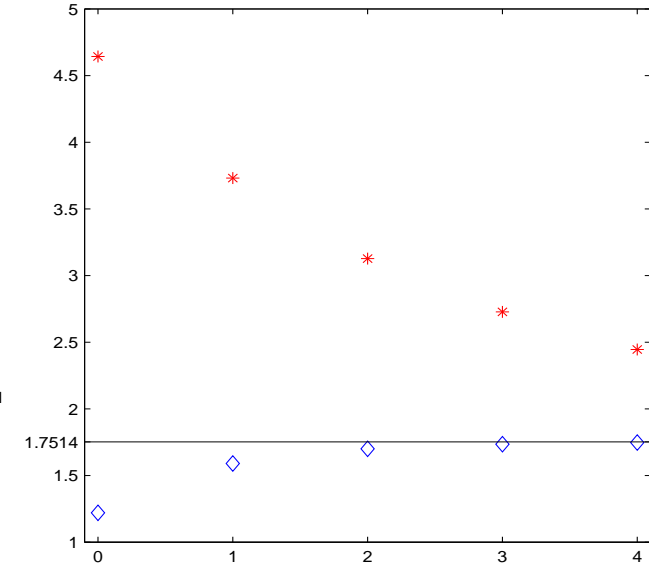
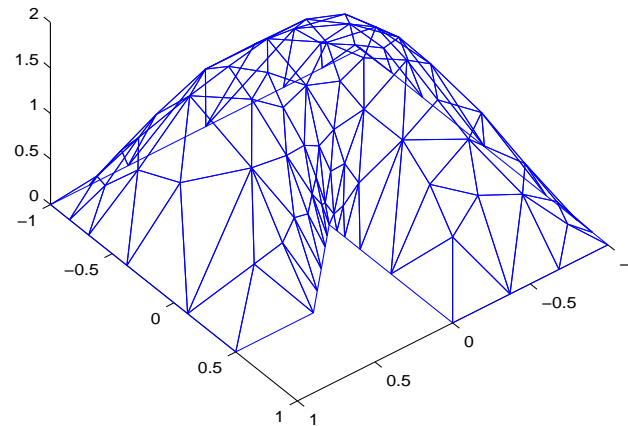
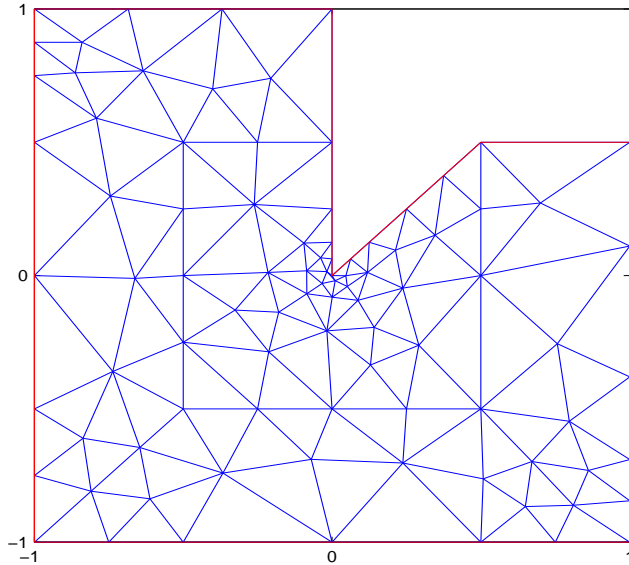
## On Computation of Lower Bounds

We usually try to have  $J(\bar{u}_{h_1}) > J(\bar{u}_{h_2}) > J(\bar{u}_{h_3}) > \dots$ , which suggests an easy way to construct the lower bounds as follows

$$a(u - u_h, u - u_h) \geq 2(J(\bar{u}_h) - J(\bar{u}_\mu)) > 0 \quad (43)$$

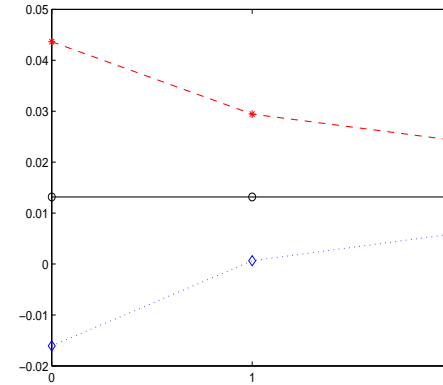
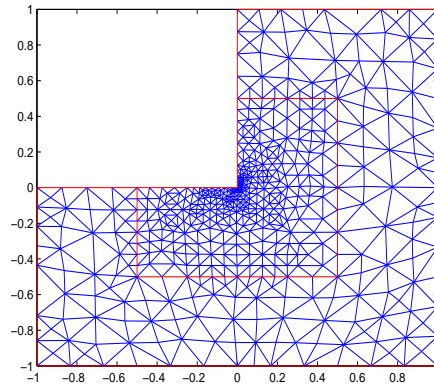
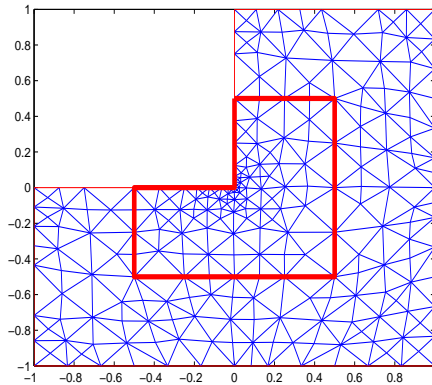
where  $\mu = h_2, h_3, \dots$

## Test No. 1



- $A \equiv I$ ,  $f \equiv 10$ ,  $\Gamma_N = \emptyset$ ,  $u_0 \equiv 0$ ,  $C_{\Omega, \Gamma_D} \leq \frac{\sqrt{2}}{\pi}$ ,  $\mathcal{T}_h$  with 92 nodes
- Several successive meshes are used in computations of bounds
- Upper bound decreases from 4.6426 to 2.4443
- Lower bound grows from 1.2191 to 1.7469
- $\|\nabla e\|^2 \approx 1.7514$  (computed using the reference solution)

## Test No. 2

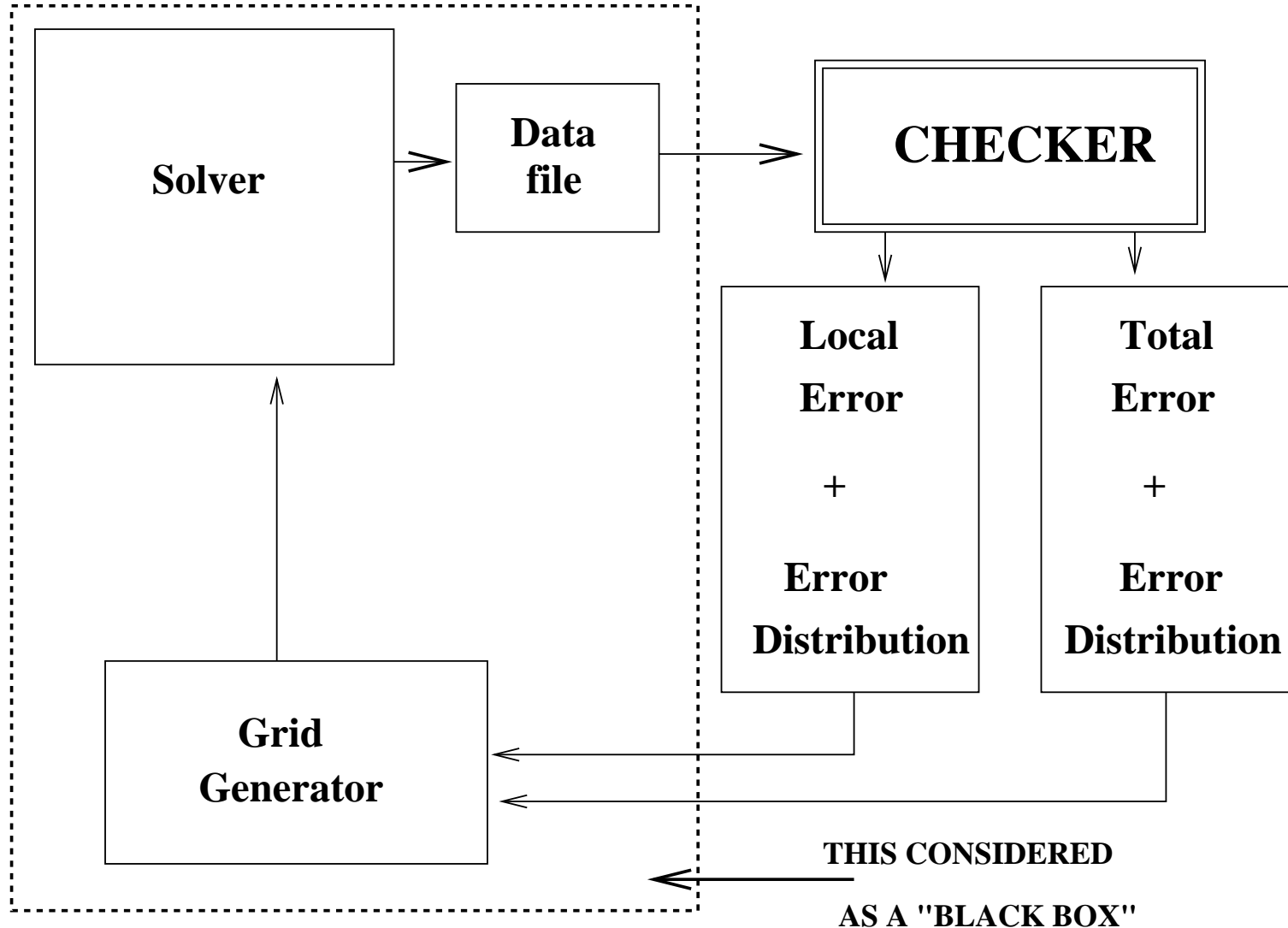


- $A \equiv I$ ,  $f \equiv 10$ ,  $\Gamma_N = \emptyset$ ,  $u_0 \equiv 0$ ,  $\mathcal{T}_h$  with 249 nodes
- $\ell(w) := \int_{\Omega} \varphi w \, dx$ ,  $\text{supp } \varphi = \omega \subset \Omega$ ,  $\varphi \in L_2(\Omega)$ ,  $\mathcal{T}_{\tau}$  with 541 nodes
- Upper bound decreases from 0.04364 to 0.02433
- Lower bound grows from -0.01606 to 0.00591
- $\ell(e) = \int_{\omega} 1 \cdot (u - u_h) \, dx \approx 0.01317$

# FINAL COMMENTS

- Techniques proposed can be adapted for the other boundary conditions and other linear elliptic problems (e.g., in linear elasticity), also for parabolic problems.
- For both types of control the corresponding upper and lower bounds can be made arbitrary close to the true errors
- Such closeness only depends on computer resources (memory, velocity) of a concrete computer used for calculations
- Technologies can be easily coded and added as **block-checker** to most of existing educational and industrial software products like **MATLAB, FEMLAB, ANSYS**, etc.

# Compatibility With Existing Codes



## Computational Benefits

By means of the above checkers, we can

- explicitly verify the error for the concrete computed solution and, therefore, see when the desired tolerance is achieved
- verify the existing codes and discover their possible drawbacks and bugs
- create optimal (or “quasi-optimal”) meshes in the process of the mesh-refinement
- improve the quality of computer-aided modelling
- reduce the time of computer simulation

- One of the main advantages of our technologies is that the error estimators proposed are valid for a wide spectrum of different approximations, i.e., the checking code is absolutely independent of the solver used
- The connection between a particular solver and the checker can be organized via the exchange of the data files



## **INPUT DATA:**

- a) geometry of the body
  - b) material constants
  - c) approximate solution obtained on the concrete mesh and the mesh parameters
- The input data are to be delivered as files of real numbers, the structure of such files is to be specified

## **OUTPUT DATA:**

- a) two-sided estimates of errors in terms of special “goal-oriented” quantities suggested by customer
  - b) two-sided bounds of the error in the global norm
  - c) distribution of local errors
- The output data can be presented both as tables of numbers and in graphical form